

# A stability theorem of Navier-Stokes flow past a rotating rigid body

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## Abstract

In this talk, I would like to consider the Navier-Stokes equations describing a viscous incompressible flow past a rotating rigid body in the 3-dimensional Euclidean space  $\mathbb{R}^3$  with axis of rotation  $\omega = a\beta_3$  ( $\beta_3 = (0, 0, 1)$ ) and with velocity  $u_\infty = k\beta_3$  where  $a$  and  $k$  are both non-zero positive constants. After transformations of coordinate and unknown functions, we have the following equations:

$$\begin{aligned} u_{tt} + u \cdot \nabla u - \Delta u + k\partial_3 u - (\omega \times x) \cdot \nabla u + \omega \times u + \nabla p &= f \\ \operatorname{div} u &= 0 \end{aligned} \tag{1}$$

in a fixed exterior domain  $\Omega$  with  $C^{1,1}$  boundary  $\Gamma$ , subject to the initial and boundary conditions:

$$\begin{aligned} u(x, t) &= \omega \times x - u_\infty && \text{on } \Gamma \\ u(x, 0) &= u_0(x) && \text{in } \Omega \\ \lim_{|x| \rightarrow \infty} u(x, t) &= 0 \end{aligned}$$

Here,  $u = (u_1, u_2, u_3)$  and  $p$  denote the unknown velocity and pressure of the fluid, respectively;  $f$  is a given external force assumed to be a function depending only on  $x$ , that is  $f = f(x)$ .

The stationary problem corresponding to equation (1) is:

$$\begin{aligned} v \cdot \nabla v - \Delta v + k\partial_3 v - (\omega \times x) \cdot \nabla v + \omega \times v + \nabla \theta &= f && \text{in } \Omega \\ \operatorname{div} v &= 0 && \text{in } \Omega \\ v(x) &= \omega \times x - u_\infty && \text{on } \Gamma \\ \lim_{|x| \rightarrow \infty} v(x) &= 0 \end{aligned} \tag{2}$$

Galdi and Silvestre [1],[2] proved a unique existence theorem of equation (2) for rather small  $k > 0$  and  $f$  in a certain norm. Moreover, they investigated a wake behind the body mathematically. Namely, their solutions have the following estimate:

$$\begin{aligned} |v(x)| &\leq M(1 + |x|)^{-1}(1 + k(|x| + x_3))^{-1} \\ |\nabla v(x)| &\leq M(1 + |x|)^{-3/2}(1 + k(|x| + x_3))^{-3/2} \end{aligned} \quad (3)$$

where  $M$  is a constant independent of  $k$  and  $a$  provided that  $k \in [0, B_1]$  and  $a \in [0, B_2]$ .

In this talk, I am interested in some stability result of the Galdi - Silvestre stationary solutions with respect to small initial disturbance, in other words the attainable problem. To explain what I want to show, I set  $u(x, t) = v(x) + w(x, t)$  and  $p(x, t) = \theta(x) + \pi(x, t)$ , where  $u$  and  $p$  are unknown functions of the original problem (1). Then,  $w$  and  $\pi$  should satisfy the following equations:

$$\begin{aligned} w_{tt} + v(x) \cdot \nabla w + w \cdot \nabla v(x) + w \cdot \nabla w \\ - \Delta w + k\partial_3 w - (\omega \times x) \cdot \nabla w + \omega \times w + \nabla \pi &= 0 \quad \text{in } \Omega \times (0, \infty) \\ \operatorname{div} w &= 0 \quad \text{in } \Omega \times (0, \infty) \\ w(x, t) &= 0 \quad \text{on } \Gamma \times (0, \infty) \\ w(x, 0) &= w_0(x) \quad \text{in } \Omega \\ \lim_{|x| \rightarrow \infty} w(x, t) &= 0 \end{aligned} \quad (4)$$

I would like to report that equation (3) admits a unique solution  $w \in C^0([0, \infty), PL_3(\Omega)) \cap C^0((0, \infty), W_3^1(\Omega) \cap L_\infty(\Omega))$  provided that  $w_0 \in PL_3(\Omega)$ ,  $\|w_0\|_{L_3}$  is small enough and  $v(x)$  and  $\nabla v(x)$  are also small enough in a certain norm, where  $P$  denotes the Helmholtz-Weyl-Leray projection. Moreover,  $w$  has the following asymptotic behaviour as  $t$  goes to infinity:

$$\begin{aligned} \|w(\cdot, t)\|_{L_q} &\leq Ct^{-\left(\frac{1}{2} - \frac{3}{2q}\right)} \quad (3 \leq q \leq \infty) \\ \|\nabla w(\cdot, t)\|_{L_3} &\leq Ct^{-\frac{1}{2}} \end{aligned}$$

When  $a = 0$  (no rotation and only past a rigid body case), in [3] I proved the same result, so that this is an extension to the case when  $a \neq 0$ .

To prove such stability result, the point is to show so called  $L_p$ - $L_q$  decay estimate of the Oseen semigroup with rotating effect. To state this more precisely, let me consider the corresponding linearized equation to (3):

$$z_t + \mathcal{L}_{k,a} z = 0 \quad (t > 0), \quad z|_{t=0} = z_0 \quad (5)$$

where  $\mathcal{L}_{k,a}$  is defined as follows:

$$\begin{aligned} \mathcal{D}_q(\mathcal{L}_{k,a}) &= \{f \in JL_q(\Omega) \cap W_q^2(\Omega) \mid (\omega \times x) \cdot \nabla f \in L_q(\Omega), \quad f|_\Gamma = 0\} \\ \mathcal{L}_{k,a} z &= P(-\Delta z + k\partial_3 z - (\omega \times x) \cdot \nabla z + \omega \times z) \quad (z \in \mathcal{D}_q(\mathcal{L}_{k,a})) \end{aligned}$$

Then, we have the following theorem.

**Theorem 1** (1) Let  $1 < q < \infty$ . Then,  $\mathcal{L}_{k,a}$  generates a  $C^0$  semigroup  $\{T(t)\}_{t \geq 0}$  on  $PL_q(\Omega)$ .

(2)  $\{T(t)\}_{t \geq 0}$  has the following estimates:

$$\|T(t)f\|_{L_r} \leq C_{q,r} t^{-\frac{3}{2}(\frac{1}{q}-\frac{1}{r})} \|f\|_{L_q} \quad t > 0$$

$$(1 \leq q \leq r \leq \infty, r \neq 1, q \neq \infty)$$

$$\|\nabla T(t)f\|_{L_r} \leq C_{q,r} t^{-\frac{1}{2}-\frac{3}{2}(\frac{1}{q}-\frac{1}{r})} \|f\|_{L_q} \quad t > 0$$

$$(1 \leq q \leq r \leq 3, r \neq 1).$$

Applying the argument in [3] based on the argument due to Kato [4] and using Theorem 1 and estimate (5), we can show the stability theorem.

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